

structure factor  $S(q)$ . However, for  $q \rightarrow 0$ , a peak develops and grows with decreasing  $T$ , again demonstrating the presence of increasing long-range behaviour for  $G_u(r, \Delta t^*)$ . No such growing peak at  $q = 0$  appears in  $S(q)$  (Fig. 4, inset). By analogy with conventional critical phenomena, we have attempted to fit  $S_u(q, \Delta t^*)$  by an Ornstein–Zernike form,  $S_u(q) \propto 1/(1 + \xi^2 q^2)$ , where  $\xi$  is the correlation length<sup>16</sup>. This form fits well at the highest  $T$ , but fails on approaching  $T_c$ , possibly because of finite size effects. Larger simulations may be required to determine accurately the correct functional form for  $S_u$  at small  $q$ . Nevertheless, the data show unambiguously that as  $T \rightarrow T_c$ , spatial correlations between the displacement of monomers become increasingly long-ranged.

Our simulations reveal a dynamical length scale relevant both to the mode-coupling dynamical transition and the glass transition. In this way we provide a bridge between the phenomenon of dynamical heterogeneity and current theories of supercooled liquids and vitrification. Remarkably, our findings are qualitatively identical to new results for simulated  $\text{Ni}_{80}\text{P}_{20}$ , a model metallic glassformer<sup>14</sup>, showing that the correlated monomer motion above  $T_c$  is neither due to nor strongly affected by chain connectivity. Instead, the striking similarity between  $\text{Ni}_{80}\text{P}_{20}$  and the present polymer melt suggests that correlated motion is a universal feature of (at least fragile<sup>1</sup>) glass-forming liquids. Furthermore, we have identified a fluctuating dynamical variable  $U$  in this polymer melt whose fluctuations become long-ranged and appear to diverge at  $T_c$ , and which thus behaves much like a static order parameter on approaching a second-order phase transition—albeit one that is not obviously accessible to traditional scattering experiments, but may be measurable in optical microscopy experiments on colloidal suspensions. Our findings suggest that substantial shifts in  $T_c$  could be observed by confining glass-forming liquids and melts, thereby limiting the divergence that can occur. Whether this could also explain the confinement-induced shifts of  $T_g$  observed experimentally<sup>22–26</sup> needs to be investigated. Our results indicate that it may be possible to obtain further insight into the nature of supercooled, glass-forming liquids using an extension to dynamically defined quantities of the framework of ordinary critical phenomena. □

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## Global changes in intensity of the Earth’s magnetic field during the past 800 kyr

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Recent advances in palaeomagnetic and dating techniques have led to increasingly precise records of the relative intensity of the Earth’s past magnetic field at numerous field sites. The compilation and analysis of these records can provide important constraints on changes in global magnetic field intensity and therefore on the Earth’s geodynamo itself. A previous compilation for the past 200 kyr integrated 17 marine records into a composite curve<sup>1</sup>, with the geomagnetic origin of the signal supported by an independent analysis of <sup>10</sup>Be production made on different cores<sup>2</sup>. The persistence of long-term features in the Earth’s magnetic intensity or the existence of long-term periodic changes cannot, however, be resolved in this relatively short time span. Here we present the integration of 33 records of relative palaeointensity<sup>3–19</sup> into a composite curve spanning the past 800 kyr. We find that the intensity of the Earth’s dipole field has experienced large-amplitude variations over this time period with pronounced intensity minima coinciding with known excursions in field direction, reflecting the emergence of non-dipole components. No stable periodicity was found in our composite record and therefore our data set does not support the hypothesis that the Earth’s orbital parameters have a direct and strong influence on the geodynamo.

We constructed the database (Table 1) by selecting records obtained after appropriate normalization of magnetization intensity. All the studies have been published with the exception of a high-resolution sequence from core ODP-1021<sup>18</sup>. In addition to the data involved in the composite curve<sup>1</sup> referred to above (‘Sint-200’), the present selection includes 16 new records relatively well distributed around the globe, 10 of which document the interval 600–800 kyr ago (Table 1). In most cases, dating was obtained by correlation to reference curves of oxygen-isotope ( $\delta^{18}\text{O}$ ) variations or other palaeoclimate proxies such as low field susceptibility or density. The age model defined by the  $\delta^{18}\text{O}$  reference curve of Bassinot *et al.*<sup>20</sup> has been applied to records older than 300 kyr; the reference curve of Martinson *et al.*<sup>21</sup> was preferred for records younger than 300 kyr because of its higher resolution. An alternative approach was used for three cores (RC10-167, KS87-752 and P226) with no obvious correlation between the palaeoclimate proxies and the  $\delta^{18}\text{O}$  reference curve. Following previous studies<sup>1,10</sup>, these individual records of palaeointensity were correlated with a preliminary stack derived from the other well dated curves. In a second

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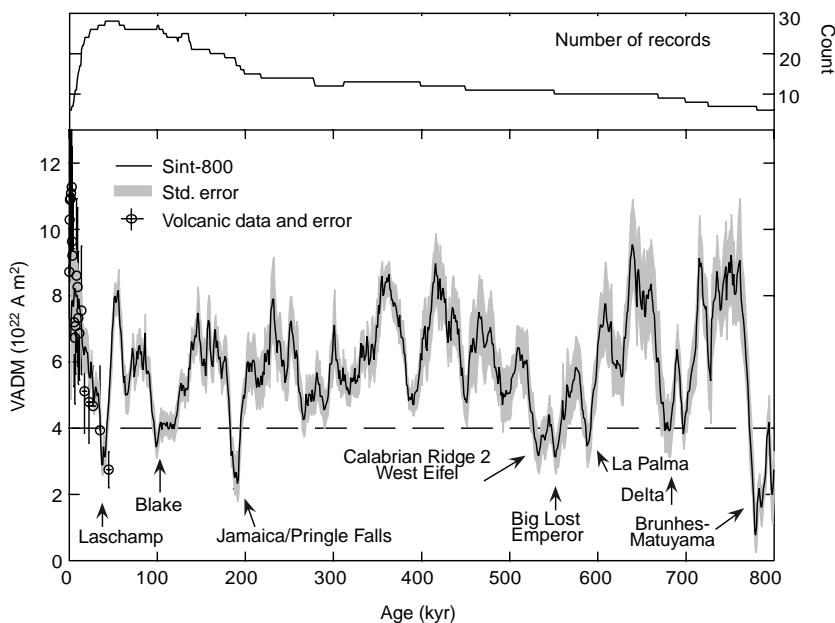
**Table 1 Records involved in the construction of Sint-800**

Record	Lat. (°N)	Long. (°E)	Sed. rate (cm kyr <sup>-1</sup> )	Normalization parameter	Other coherent normalization	Dating	Interval used (kyr)	Ref.
Set no. 1								
KET82-51	39	14	6.2	χ	ARM, SIRM	Corr. with KET8004 + tephra.	10-90	3
MD84-629	36	33	11.6	χ	ARM, SIRM	δ <sup>18</sup> O + tephra.	20-60	3
DED87-707	40	13	10.4	χ	ARM, SIRM	Corr. with KET8004 + tephra.	10-60	3
MD85-668	-1	46	6.2	ARM	χ, IRM	δ <sup>18</sup> O	20-140	4
MD85-669	2	47	5.3	ARM	χ, IRM	χ corr. with MD668	20-140	4
MD85-674	3	50	10.9	ARM	χ, IRM	χ corr. with MD668	20-140	4
ERDC113P	2	159	1.2	χ		δ <sup>18</sup> O	10-360	5, 6
NGC29	4	136	2.4	ARM	χ	χ correlated with NP5	0-190	7
NGC26	2	135	3.7	ARM	χ	χ correlated with NP5	0-120	7
NGC16	2	135	3.7	ARM	χ	χ correlated with NP5	0-190	7
NP7	1	138	2.0	ARM	χ	χ correlated with NP5	10-200	7
NP5	1	136	3.7	ARM	χ	δ <sup>18</sup> O	10-200	7
SU92-18	37	-27	3.5	ARM	χ, SIRM	δ <sup>18</sup> O	0-280	8
SU92-19	37	-27	3.2	ARM	χ, SIRM	Reflectance + χ correlated with SU92-18	0-280	8
ODP-768A	8	121	9	χ	ARM, IRM	δ <sup>18</sup> O + AMS radiocarb.	10-90	9
ODP-768B	8	121	10	χ	ARM, IRM	δ <sup>18</sup> O + AMS radiocarb.	10-130	9
ODP-769B	8	121	8.2	χ	ARM, IRM	δ <sup>18</sup> O + AMS radiocarb.	10-160	9
P012	58	-47	10	IRM	χ, ARM	δ <sup>18</sup> O + AMS radiocarb. + calib. Sint-200	10-180	10
P013	58	-48	10-30	IRM	χ, ARM	δ <sup>18</sup> O + AMS radiocarb. + calib. Sint-200	10-220	10
P094	50	-45	10	IRM	χ, ARM	δ <sup>18</sup> O + AMS radiocarb. + calib. Sint-200	0-110	10
Set no. 2								
ODP-983	60	-24	10.4	IRM	χ, ARM	δ <sup>18</sup> O	0-725	11
ODP-984	61	-24	12.6	IRM	χ, ARM	δ <sup>18</sup> O	0-450	11
KS87-752	-38	-38	2.4	ARM	χ	χ corr. with ODP940	310-800	12
ODP769	9	121	13	χ	ARM, IRM	δ <sup>18</sup> O	200-800	13
NP35	4	141	1	ARM	IRM	δ <sup>18</sup> O	130-700	14
RNDB75P	2	160	1.5	IRM	χ, ARM	δ <sup>18</sup> O	120-670	6
MD940	-6	61	1.3	ARM	χ, IRM	χ corr. with ODP709	100-800	15
ODP851	2	-110	1.8	ARM	χ, IRM	GRAPE, astro. calib.	30-800	16, 17
P226	3	-169	0.6	ARM	IRM	B/M reversal	40-780	14
ODP1021	40	-128	3.7	χ	ARM	χ, astro. calib.	10-800	18
NGC38	-15	-175	1	ARM	IRM	δ <sup>18</sup> O	10-400	14
RC10-167	33	-150	2.1	ARM		B/M reversal	0-800	19
NGC36	1	160	1	ARM	IRM	δ <sup>18</sup> O	0-550	14

χ, magnetic susceptibility; ARM, anhysteretic remanent magnetization; IRM, isothermal remanent magnetization. Data available as Supplementary Information.

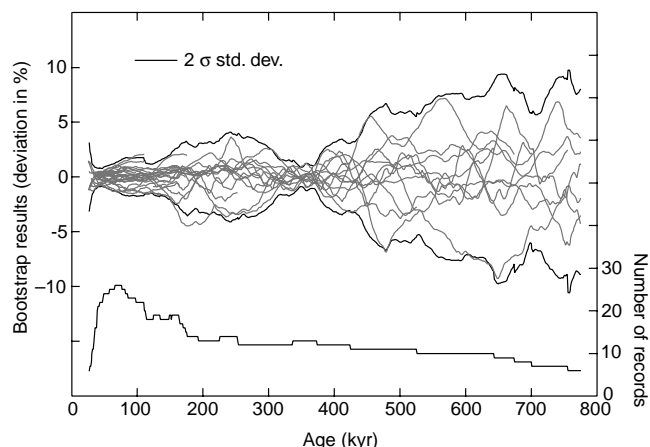
step, the age models were applied to the susceptibility records and subsequently refined with respect to the δ<sup>18</sup>O reference curve. This procedure increased the quality of the correlation by a factor of two. All the records have been linearly interpolated every 1 kyr before calculation of the stack.

Palaeointensity variations can be extracted from natural remanent magnetization (NRM) after normalization of the intensity by an appropriate magnetic parameter which activates the same magnetic fraction as that which carries the NRM. However, various magnetic parameters and different magnetization levels have been



**Figure 1** Synthetic record (Sint-800) with its standard error obtained from the stack of 33 records of palaeointensity. The upper part of the figure gives the number of records within successive time intervals. All the records have been dated with a common timescale and normalized (see text). The curve has been calibrated over the past 40 kyr in terms of virtual axial dipole moments (VADMs)

using averaged volcanic data (open circles) over 5-kyr intervals (correlation coefficient,  $r = 0.7$ ). The horizontal dashed line corresponds to the critical value of intensity below which directional excursions have been observed (Langereis *et al.*<sup>25</sup>). Other dips of lower amplitude coincide also with excursions reported from marine sediments.

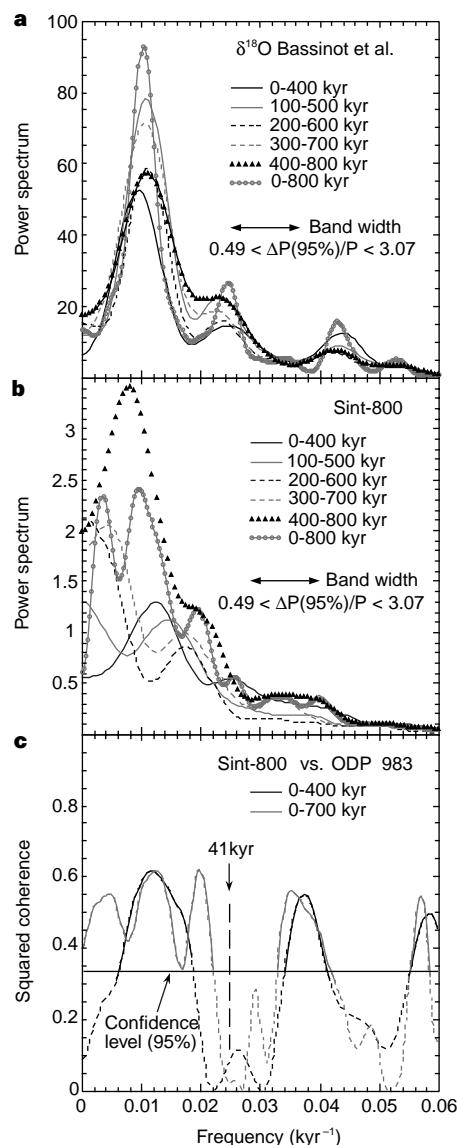


**Figure 2** Tests of our synthetic record, Sint-800. Shown are results of a bootstrap process aimed at identifying drawings of 33 different composite records that show significant deviations from Sint-800. Each plot corresponds to the divergence between Sint-800 and another composite calculated after removing one record from the database. No outlier can be identified.

used by authors. Integration of the records on a common scale can therefore only be done after division of the intensities by their mean value over the entire time interval. This simple procedure accounts also for the latitudinal dependence of the dipole field intensity. As the average geomagnetic field was probably different over different periods, this normalization should be done over the same time interval for all the records. Unfortunately, no time interval is common to the entire data set which was thus split into two different groups. Records from the first group (Table 1) are younger than 350 kyr and they were normalized to unity over the interval 24–56 kyr ago. The interval 312–398 kyr ago was used for the second group. Subsequently, two independent stacks were constructed and compared over the common interval 0–350 kyr. Final integration of the two data sets was obtained after correcting the records of the first group by a factor 0.84, which corresponds to the ratio between the two averaged values.

In Fig. 1 we show our composite record ‘Sint-800’ and the associated standard error obtained after stacking all the data. In the same figure are plotted also the number of records involved over each successive time interval. The stability of the results has been tested by a bootstrap technique similar to that used for Sint-200. No significant discrepancy was observed between the stacks performed from 33 successive calculations. Deviations between Sint-800 and each successive drawing (after smoothing with a 50-kyr-long moving window to slightly reduce the resolution) lie within the limits imposed by the  $2\sigma$  standard deviation (Fig. 2) and confirm the absence of significant outliers. We also note that the sharpness of the distribution is related to the number of records. Thus, some refinement (mainly a reduction of the uncertainties) would be produced by incorporation of future records covering the interval 500–800 kyr ago. Conversion into virtual axial dipole moments (VADMs) were done using volcanic records of absolute palaeointensity. Volcanic data must be averaged out over finite time intervals in order to eliminate the contribution of the non-dipole field. Only the period 0–40 kyr ago<sup>22</sup> is sufficiently documented to allow this. According to this calibration, the highest intensities culminate at about  $9 \times 10^{22}$  A m<sup>2</sup> and the time-averaged field value during the Brunhes chron was  $(6.0 \pm 1.5) \times 10^{22}$  A m<sup>2</sup>.

One characteristic of Sint-800 (Fig. 1) is that many intensity lows occurred during the past 800 kyr. Sint-800 was constructed from world-wide marine records with a resolution that does not exceed 3–5 kyr and is thus sensitive to the variations of the axial dipole moment. A simple condition for the occurrence of excursions<sup>23,24</sup> is that the dipole field remains abnormally low so that non-dipole



**Figure 3** Spectral analysis of Sint-800 using the Blackman–Tukey method with a Bartlett window. **a, b**, Power spectra of the  $\delta^{18}\text{O}$  curve<sup>20</sup> (**a**) and Sint-800 (**b**). Analyses have been performed over different time intervals (to explore the stability of the results). There is no stationarity or periodicity emerging from the palaeointensity variations. **c**, Coherence function between Sint-800 and ODP Site 983 (ref. 11). The calculations involve the entire 0–700 kyr interval common to the two curves and then only the 0–400 kyr interval which incorporates more records.

effects become dominant. We referred to the analysis of the short events by Langereis *et al.*<sup>25</sup> to select all the excursions that have been previously reported from both sedimentary and volcanic records. To those we added the discovery<sup>26</sup> of a large field excursion data at  $595 \pm 14$  kyr ago from a volcanic sequence at La Palma. All these events coincide with periods of significant intensity drops (Fig. 1) and occur below a critical field of  $\sim 4 \times 10^{22}$  A m<sup>2</sup>, that is, less than 50% of the present-day field. In fact, a reduction of the present-day dipole intensity by this amount would induce the emergence of non-dipole components at many locations<sup>25,27</sup>. Other excursions reported from marine sediments<sup>23,25,27</sup> are synchronous with dips of lower amplitude. There is no correlation between these intensity dips and cold climate events, although such a correlation has been suggested<sup>28</sup>.

The possibility that the geomagnetic field variations would be modulated by the Earth’s orbital parameters has been suggested by

several authors<sup>29–32</sup>. Nevertheless, most studies dealing with long and well dated sequences conclude either that a 30–50 kyr geomagnetic pulse<sup>3,6</sup> exists, or that there is no dominant period. But a correlation has been found<sup>11</sup> between the 41-kyr obliquity cycle and palaeointensity variations; this was based on the analysis of a very detailed and precisely dated sequence<sup>11</sup> from Ocean Drilling Project Site 983. We performed a spectral analysis of Sint-800 using the Blackman–Tukey technique with the AnalySeries software<sup>33</sup>. In order to investigate the stability of the results, we first analysed the entire signal and then restrained the analysis to the 0–400 kyr interval which incorporates more records. The spectral contents of these two intervals are significantly different. The same conclusion is reached by comparing every 400-kyr-long interval in increments of 100 kyr with the 0–800 kyr period (Fig. 3). These results point out the absence of any dominant stable periodicity. For comparison, the same analysis performed with the  $\delta^{18}\text{O}$  curve<sup>20</sup> shows perfect reproducibility of the orbital peaks (23, 41 and 100 kyr) over each interval. (An additional indication supporting these conclusions is that the artificial 41-kyr signal created after band-pass-filtering Sint-800 has a different phase and a much smaller amplitude than the signal obtained by the same procedure for site 983.) Similar observations and conclusions are reached without incorporating the site 983 record in the database. Orbital modulation of the geomagnetic field should primarily affect the dipole field, and therefore should be present in Sint-800 which incorporates at least 20 data points per 41 kyr. However, for periods shorter than 40 kyr it is possible that uncertainties due to dating and/or smoothing inherent to the stacking process induced loss of spectral power (Fig. 3b). This may explain the pattern observed in the coherence functions between Sint-800 and data from site 983 (Fig. 3c) for these short periods.

Our composite Sint-800 curve shows that the Earth's dipole field over the past 800 kyr was dominated by changes of very large (and various) amplitude but does not indicate the presence of any dominant periodicity. The mean field value remained more or less constant. Geomagnetic excursions are observed when the dipole moment decreases to a critical value of about  $4 \times 10^{22} \text{ A m}^2$ , and such excursions must thus be seen as direct consequences of the overall 'secular' variation of the dipole field. □

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Supplementary information is available on Nature's World-Wide Web site (<http://www.nature.com>) or as paper copy from the London editorial office of Nature.

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## Nature of the Earth's earliest crust from hafnium isotopes in single detrital zircons

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Continental crust forms from, and thus chemically depletes, the Earth's mantle. Evidence that the Earth's mantle was already chemically depleted by melting before the formation of today's oldest surviving crust has been presented in the form of Sm–Nd isotope studies of 3.8–4.0 billion years old rocks from Greenland<sup>1–5</sup> and Canada<sup>5–7</sup>. But this interpretation has been questioned because of the possibility that subsequent perturbations may have re-equilibrated the neodymium-isotope compositions of these rocks<sup>8</sup>. Independent and more robust evidence for the origin of the earliest crust and depletion of the Archaean mantle can potentially be provided by hafnium-isotope compositions of zircon, a mineral whose age can be precisely determined by U–Pb dating, and which can survive metamorphisms<sup>4</sup>. But the amounts of hafnium in single zircon grains are too small for the isotopic composition to be precisely analysed by conventional methods.